Magnetic Structure and Glassiness in Fe$_{0.5}$Ni$_{0.5}$PS$_3$

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Abstract

This work explores the magnetic properties of Fe$_{0.5}$Ni$_{0.5}$PS$_3$. The system shows pronounced hysteresis in the magnetic phase transition temperature as a function of the direction of the change in temperature. Field cooled/zero field cooled hysteresis is not pronounced. However, the transition temperature between antiferromagnetic and paramagnetic order occurs at approximately 97 K on cooling, but at 138 K on warming, whether the warming is after zero field or field cooling.

This is indicative of magnetic glassiness, and made all the more unusual because all measurements exhibit a transition to a third magnetic phase existing at temperatures below ~14 K. The intermediate phase relaxes on a laboratory time scale of the order of 48 minutes, into an antiferromagnetic state whose magnetic structure is, from neutron diffraction, indistinguishable from the low temperature state. This low temperature state shows magnetic ordering consistent with that observed in CoPS$_3$ and NiPS$_3$. Analysis of the neutron measurements shows that the direction of moments cannot be along the b axis. It is also shown that the moments are unlikely to lie in the c* direction. Therefore, we suggest that the moments lie along the a axis.

Keywords: antiferromagnetism, layered magnetism, FePS$_3$, magnetic glass
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1. Introduction

The family of two-dimensional magnetic materials MPS$_3$ where $M = Mn^{2+}$, Fe$^{2+}$, Ni$^{2+}$, Zn$^{2+}$ etc shows a wide range of magnetic behaviour [1, 2, 3, 4, 5, 6, 7]. Given that Fe$_{0.5}$Mn$_{0.5}$PS$_3$ has been found to be a spin glass [2], it was
considered to be of interest to explore the behaviour of Fe$_{0.5}$Ni$_{0.5}$PS$_3$. The end members, FePS$_3$ and NiPS$_3$, are antiferromagnets with transition temperatures of 123 K and 155 K respectively.

In FePS$_3$ and NiPS$_3$ the correlation structure of the spins was thought to be similar [8], with the nature of the anisotropy different. However, more recent work suggests that the magnetic structure of FePS$_3$ is not yet fully understood [9]. The competing anisotropies – XY in NiPS$_3$ and Ising in FePS$_3$ – are also likely to lead to interesting magnetic behaviour. The three-fold in-plane coordination in these honeycomb lattice systems renders geometrical frustration unlikely, despite the significant next-nearest neighbour interactions (for example [10]).

The crystal structure of a typical member of the MPS$_3$ family, MnPS$_3$, is shown in figure 1 and shows the honeycomb arrangement of Mn atoms, the P$_2$ dimers and the layers of sulphur atoms that sandwich the Mn/P$_2$ layers. There is a van der Waals gap between these sandwich structures resulting in an indirect exchange path along c, contributing to the 2D magnetic properties and giving the crystals highly anisotropic structural behaviour.

The crystal structure of manganese disulfide (MnS$_2$) and the space group is C$\overline{2}/m$ and lattice parameters are $a = 5.812(2)$ Å, $b = 10.070(3)$ Å, $c = 6.632(1)$ Å and $\beta = 106.98(3)^\circ$ (NiPS$_3$, [11]) and $a = 5.934(6)$ Å, $b = 10.28(1)$ Å, $c = 6.772(7)$ Å and $\beta = 107.2(1)^\circ$ (FePS$_3$, [12]).

A magnetic glass generally occurs occurs when a transition between magnetic states arrested through some mechanism, such as competing interactions and/or anisotropy. This gives rise to a state consisting of randomly oriented clusters of moments, frozen over time scales of minutes/hours [14]. Three measurements that can be used to distinguish spin glasses and magnetic glasses are:

- Thermal hysteresis: For a magnetic glass a distinct thermal hysteresis occurs between the FCC and FCW susceptibilities (FCC is field cooled
cooling — measurement made in a field while cooling the sample and FCW is field cooled warming — measurements made while warming in a field after cooling in a field. For spin glasses, the FCC and FCW curves should be identical.

- Irreversibilities: The thermomagnetic irreversibility of a magnetic glass appears above a critical field (dependent on the material) and increases with increasing applied field. The thermomagnetic irreversibility of a spin glass decreases with increasing applied field.

- Relaxation: The FCC/FCW states of the magnetic glass undergo relaxation towards the equilibrium ZFC state, regardless of applied field.

2. Experimental

2.1. Sample Preparation

Samples of FePS$_3$, NiPS$_3$ and Fe$_{0.5}$Ni$_{0.5}$PS$_3$ were synthesised by direct combination of the appropriate metal sulphides with phosphorus and sulphur. Stoichiometric weights of the powders were thoroughly homogenised using a mortar and pestle, pressed into pellets and sealed inside quartz tubes which were flushed with argon and evacuated to $10^{-3}$ Torr. To avoid the build up of pressure inside the ampoule, the FePS$_3$ and NiPS$_3$ samples were ramped to 700°C over 8 days, before being held at 700°C for 7 and 11 days respectively. The Fe$_{0.5}$Ni$_{0.5}$PS$_3$ sample was ramped over 24 days due to larger sample mass, before being held at 700°C for 1 month. It was found that regrinding and a second anneal was required to achieve homogeneity in mixed metal samples such as Fe$_{0.5}$Ni$_{0.5}$PS$_3$.

2.2. X-ray Diffraction

Powder XRD data were collected on a Siemens D5000 Diffractometer. A monochromator in front of the detector selected out Cu Kα radiation and removed the X-ray fluorescence originating from iron in the sample. Rietica [15] was used to perform Rietveld [16] and Le Bail [17] fits.

2.3. Electron Microscopy

Compositional analysis was undertaken using a Joel 6400 Scanning Electron Microscope (SEM) at the ANU Centre for Advanced Microscopy to determine the composition of Fe$_{0.5}$Ni$_{0.5}$PS$_3$. Samples of the end members were also analysed for use as standards.

2.4. Magnetisation Measurements

Magnetic susceptibility measurements used a 1.6 T EG&G PARC model number 155 Vibrating Sample Magnetometer (VSM) and closed-cycle helium refrigerator (CCR) at UNSW Canberra Campus. In all experiments the sample was first cooled to 8.0K in zero applied field. To identify hysteretic magnetic behaviour three different measurement protocols were used.
1. Zero Field Cooled Warming (ZFCW): After initial cooling without a field, a field of 1T was applied and the sample was warmed at a constant rate of 3K/min to 320K.

2. Field Cooled Cooling (FCC): The field was maintained and the sample was cooled to 8.0K. Due to the nature of the CCR, the cooling rate could not be as tightly controlled as the warming rate, but over the most significant span (200K to 75K) was reasonably consistent at -7K/min.

3. Field Cooled Warming (FCW): The sample was maintained at 8K for approximately an hour and then field was maintained while the sample was warmed back to room temperature at a constant rate of 3K/min.

To confirm magnetic relaxation of Fe$_{0.5}$Ni$_{0.5}$PS$_3$, the ZFCW and FCC protocols were repeated, however the sample was then cooled to 100K during the FCC cycle and held at this temperature for 3 hours. The sample was a loosely packed powder.

These experiments were repeated (with identical temperature ranges and equivalent cooling/warming rates) with the field applied (1) parallel and (2) orthogonal to the compression direction of a small uniaxially pressed pellet of Fe$_{0.5}$Ni$_{0.5}$PS$_3$.

2.5. Neutron Diffraction

Neutron powder diffraction data were collected at 1.62158 Å using the Echidna diffractometer at the OPAL Reactor (Bragg Institute, Australian Nuclear Science and Technology Organisation) [18]. Measurements were made at a range of temperatures in magnetic fields of 0T and 1T.

Because of the plate-like shape of the crystallites of Fe$_{0.5}$Ni$_{0.5}$PS$_3$, the sample consisted of three uniaxially pressed 6mm diameter pellets mounted with mutually orthogonal axes. The sample could not be rotated during the experiment. This arrangement ensured that more of reciprocal space was sampled than would have been the case using a single pellet with its strong preferred orientation (so strong that some classes of reflections may be suppressed entirely). However, each of the three pellets had strong preferred orientation, meaning that while the diffraction pattern was more thoroughly sampled, relative peak intensities were not reliable. This means that it was possible to identify and index magnetic reflections of all classes (something that would have been impossible in a highly oriented sample), which allowed determination of the correlation structure of the magnetic moment. However, only limited information could be gained regarding the magnetic moment magnitudes and directions.

The measurements were made while the sample was in its ‘relaxed’ state to ensure no transient effects were observed. This was achieved by allowing the sample to equilibrate for approximately two hours between successive measurements. This was necessary as the data collection times were long relative to the sample relaxation time, and measurements made while the sample was relaxing would be uninterpretable.
### Table 1: Lattice parameters for FePS₃ (Rietveld refinement), NiPS₃ and Fe₀.₅Ni₀.₅PS₃ (Le Bail refinement) in the monoclinic space group C2/m.

<table>
<thead>
<tr>
<th>Sample</th>
<th>a (Å)</th>
<th>b (Å)</th>
<th>c (Å)</th>
<th>β (°)</th>
<th>V (Å³)</th>
</tr>
</thead>
<tbody>
<tr>
<td>FePS₃</td>
<td>5.952(1)</td>
<td>10.305(1)</td>
<td>6.751(1)</td>
<td>107.39(1)</td>
<td>395.1</td>
</tr>
<tr>
<td>NiPS₃</td>
<td>5.817(1)</td>
<td>10.082(1)</td>
<td>6.627(1)</td>
<td>106.91(1)</td>
<td>371.8</td>
</tr>
<tr>
<td>Fe₀.₅Ni₀.₅PS₃</td>
<td>5.893(5)</td>
<td>10.193(5)</td>
<td>6.651(5)</td>
<td>106.76(5)</td>
<td>382.5</td>
</tr>
</tbody>
</table>

### Table 2: Atomic coordinates in Fe₀.₅Ni₀.₅PS₃; occupancy and ADPs not refined.

<table>
<thead>
<tr>
<th>Atom</th>
<th>x</th>
<th>y</th>
<th>z</th>
<th>Occupancy, n (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Fe</td>
<td>0</td>
<td>0.327(2)</td>
<td>0</td>
<td>50</td>
</tr>
<tr>
<td>Ni</td>
<td>0</td>
<td>0.327(2)</td>
<td>0</td>
<td>50</td>
</tr>
<tr>
<td>P</td>
<td>0.040(4)</td>
<td>0</td>
<td>0.121(2)</td>
<td>100</td>
</tr>
<tr>
<td>S₁</td>
<td>0.730(3)</td>
<td>0</td>
<td>0.239(2)</td>
<td>100</td>
</tr>
<tr>
<td>S₂</td>
<td>-0.263(2)</td>
<td>0.336(1)</td>
<td>0.244(1)</td>
<td>100</td>
</tr>
</tbody>
</table>

### 3. Crystal Structure

Crystal structures of samples of FePS₃ and NiPS₃ were established to be close to those expected [11, 12]. Parameters were established for Fe₀.₅Ni₀.₅PS₃. Table 1 lists key unit cell parameters. Figure 2 shows the XRD pattern for Fe₀.₅Ni₀.₅PS₃, along with a Rietveld fit. Table 2 gives atomic coordinates; atomic displacement parameters (ADPs) and occupancy factors could not be refined reliably due to the strong preferred orientation.

The cell volume of Fe₀.₅Ni₀.₅PS₃ is close to the average of the volumes of FePS₃ and NiPS₃. Based on Vegard’s law, the experimentally achieved composition FeₓNi₁₋ₓPS₃ is x = 0.50 ± 0.03. SEM-EDX, averaged over a range of crystallites in the powder, gives x = 0.52 ± 0.03, indicating a sample very close to the desired composition.

### 4. Magnetisation Measurements

The magnetisation of powder samples of FePS₃ and NiPS₃ were measured as functions of temperature after cooling samples in zero field. Samples were heated in applied fields of between 100 mT and 1 T from 4 K to room temperature. The observed transition temperatures agreed with those from the literature. Because of the powder nature of the samples, the parallel and perpendicular susceptibilities could not be separated, and the sample susceptibility, χ, was a combination of the two. The platiness of the crystal habit means it is unlikely that the observations reflect a simple average of χ = 1/3χ∥ + 2/3χ⊥.

Figure 3 shows the magnetisation measurements for a loosely packed powder of Fe₀.₅Ni₀.₅PS₃ at 1 T; measurements made at 100 mT are similar but noisier. Equivalent measurements were made on FePS₃ and NiPS₃ and the hysteresis apparent in Fe₀.₅Ni₀.₅PS₃ was not seen. Measurements performed on uniaxially
pressed pellets of Fe$_{0.5}$Ni$_{0.5}$PS$_3$ showed hysteresis, suggesting that it is not related to, for example, grain rotation in the applied field. Similarly, the difference between $T^2_N$ on cooling and warming was not a result of poor heat transfer during the experiment. The results were repeated for independently manufactured and measured samples.

We note that there is some noise in the signal, the result of measuring relatively small antiferromagnetic moments on an instrument more suited to systems with ferromagnetic moments.

The sample shows three magnetic phases — paramagnetism above the cusp at around 97K (this temperature is denoted $T^2_N$), and antiferromagnetism below the lower cusp at approximately 13K ($T^1_N$), and an intermediate phase that also appears antiferromagnetic. This is essentially in line with what has been previously observed [19, 5], where the two transitions were assigned to different sublattices. However, the earlier work undertook ZFCW and FCW measurements and saw no difference, and concluded that the system was not glassy, as indeed it is not a conventional spin glass. However, figure 3 shows that there is pronounced hysteresis as a function of the direction of the change in temperature, that is, FCC/FCW (or FCC/ZFCW) hysteresis, as distinct from ZFCW/FCW hysteresis as observed in a conventional spin glass. This is one signature of a ‘magnetic glass’. For conventional spin glasses, the FCC and FCW curves should be identical [14].

While [5] notes that “Magnetic susceptibility measurements were done from 300K down to 10K” it can be noted that ZFC measurements must be done while warming after cooling in zero field, and thus we do not interpret this comment as meaning that measurements we made on cooling in [5], and thus there is no...
Figure 3: Magnetic susceptibility of a powder of Fe$_{0.5}$Ni$_{0.5}$PS$_3$, measured in a field of 1 T under a range of conditions. Arrows indicate temperature increasing or decreasing, while the order in which measurements were made was: (1) warming after cooling in zero field (ZFCW); (2) cooling in a field (FCC) and lastly; (3) warming after cooling in a field (FCW).

discrepancy between that work and results presented here.

Curie-Weiss fits to these (admittedly noisy) data for $T > T_N^2$ give an effective ordered moment of $2.86 \pm 0.06\mu_B$ for all measurements, while for the warming runs $\theta = 77 \pm 1$K and for the cooling run there is a small difference and $\theta = 73 \pm 1$K. $2.86\mu_B$ is very close to the quenched moment for Ni$^{2+}$ given by Hund’s rules, whereas an average of Fe$^{2+}$ and Ni$^{2+}$ quenched moments gives $3.87\mu_B$, as observed in [19]. It should be noted however that [19] also suggest different critical temperatures for the Ni and Fe sublattices, or at least different critical temperatures ‘related to’ the sublattices. If this is taken to mean transitions to paramagnetism happen separately on the two sublattices, this would imply different values of $\theta$ for Fe versus Ni in a Curie-Weiss fit, as the Fe would be making a paramagnetic contribution from a much lower temperature. It is not clear that this has been catered for in their analysis. If the assumption is made that both Ni and Fe show their quenched moments in the paramagnetic regime, but that Fe ‘disorders’ at a much lower temperature, it is possible to fit the data in figure 3 to an expression of the form

$$\chi = \frac{C_{Ni}}{T + \theta_{Ni}} + \frac{C_{Fe}}{T + \theta_{Fe}}$$  \hspace{1cm} (1)
where the usual Curie-Weiss law for an antiferromagnet is $\chi = \frac{2C}{T + \theta}$ and half the magnetic atoms are Fe and half Ni. Taking the moments of Ni and Fe as 2.83$\mu_B$ and 4.90$\mu_B$ respectively, and fitting only the $\theta$ values in equation 1 it is possible to get identically ‘good’ fits to the data as fitting a single Curie-Weiss function in which both $C$ and $\theta$ are varied. Hence it is difficult to conclude whether the model in which there are separate sublattices for Ni and Fe is valid.

To gain some insight into the directions of the moments, measurements were made on a uniaxially pressed pellet, with the compression axis parallel and then perpendicular to the applied field. If the moments were close to $c^*$ (as they are in MnPS$_3$, where they are at an angle of 8$^\circ$ to $c^*$ [20]) then this should induce a substantial difference between susceptibility measured for the field parallel to the compression direction compared to a perpendicular measurement. However, if the moments are in the plane differences would be less pronounced. As no substantial difference between susceptibility was observed for the two orientations it was determined that the magnetic moments are most likely not aligned along $c^*$.

Once the first condition for magnetic glassiness was established — the presence of hysteresis — (see section 1) — a second test was performed to test for relaxation in the system. The sample was cooled in zero field, measured while warming in a field, giving a conventional ZFCW measurement, and the the FCC measurement was begun, but the cooling was halted at 100K — below $T_N$ as observed on warming but above $T_N$ as observed on cooling. The result is shown in figure 4.

The relaxation was fitted to a stretched exponential for the form

$$\chi = A + Be^{-(t/t_0)^\beta}$$  \hspace{1cm} (2)

where $A$ and $B$ scale the fit to the data, $t$ is time, $t_0$ is the ‘start time’ for the relaxation, $\tau$ is the time constant and $\beta$ is the stretching parameter. The effective mean relation time constant, $\langle \tau \rangle$, is given by

$$\langle \tau \rangle = \frac{\tau}{\beta \Gamma \left( \frac{1}{\beta} \right)}$$  \hspace{1cm} (3)

where $\Gamma$ is the gamma function.

$\tau = 12.5(5)\text{sec}$ and $\beta = 0.186(6)$ gives $\langle \tau \rangle \sim 48$ minutes. The fit is extremely good, and suggests that the stretched exponential, which is often used for relaxational phenomena in disordered systems, is appropriate here.

The final value of $\chi$ is within error of that seen in the ZFCW experiment, as would be expected for a magnetic glass. This time constant shows the system is far from completely frozen, yet relaxes slowly compared to typical time scales for magnetic susceptibility measurements as a function of $T$.

5. Neutron Diffraction

To identify magnetic diffraction peaks neutron diffraction patterns were collected at 0K, 70K, 100K and 225K. Purely magnetic diffraction peaks show zero
Figure 4: Relaxation of magnetic susceptibility of a powder of Fe$_{0.5}$Ni$_{0.5}$PS$_3$, measured in a field of 1 T at 100K. Decrease of susceptibility from FCC value to average of ZFCW/FCW values over a period of approximately 2 hours.
Figure 5: Low angle region of three diffraction patterns measured with 1.62158 Å neutrons. Magnetic and nuclear reflections identified from differences between traces — 225K (lower trace, light grey) 100K (middle trace, dark grey) and 5K (upper trace, black) ‘N’ indicates a nuclear reflection, and ‘M’ a magnetic. The 011 is very weak. The inset shows the temperature dependence of the integrated intensity of the 010 peak; the point noted ‘VSM’ gives $T_N^2$ as given by magnetometry (ZFCW).

normalised intensity at 225K as $\text{Fe}_{0.5}\text{Ni}_{0.5}\text{PS}_3$ is paramagnetic at this temperature.

The magnetic reflections obey $h + k = 2n + 1$ which is not allowed for the $C2/m$ space group associated with the crystal structure of $\text{Fe}_{0.5}\text{Ni}_{0.5}\text{PS}_3$. These reflections agree with those for the CoPS$_3$ and NiPS$_3$ magnetic structures [21]. The difference between the CoPS$_3$ and NiPS$_3$ magnetic structures is the direction of spin, namely along the $a$ axis and $c$ axis respectively.

As the 010 magnetic reflection is present in $\text{Fe}_{0.5}\text{Ni}_{0.5}\text{PS}_3$ at low $T$, the direction of spin of the magnetic moments cannot lie along the $b$ axis — magnetic reflections with scattering vector parallel to the moment direction show no magnetic structure factor. Results from section 4 suggest that the moments are unlikely to be aligned along $c^*$ (or $c$, which is close to but not parallel with $c^*$). Therefore, the moments are likely to be either aligned along the $a$ axis or an intermediate direction. Preferred orientation prevents the determination of moment magnitude.
Table 3: Lattice parameters for Fe\textsubscript{0.5}Ni\textsubscript{0.5}PS\textsubscript{3} (Le Bail refinement) in the monoclinic space group C\textsubscript{2}/m, with T.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>RT (XRD)</th>
<th>225K (neutron)</th>
<th>100K (neutron)</th>
<th>70K (neutron)</th>
<th>5K (neutron)</th>
</tr>
</thead>
<tbody>
<tr>
<td>a (Å)</td>
<td>5.893(5)</td>
<td>5.887(3)</td>
<td>5.863(3)</td>
<td>5.863(2)</td>
<td>5.860(2)</td>
</tr>
<tr>
<td>b (Å)</td>
<td>10.193(9)</td>
<td>10.168(5)</td>
<td>10.183(5)</td>
<td>10.182(4)</td>
<td>10.178(3)</td>
</tr>
<tr>
<td>c (Å)</td>
<td>6.651(6)</td>
<td>6.668(5)</td>
<td>6.654(6)</td>
<td>6.651(5)</td>
<td>6.647(4)</td>
</tr>
<tr>
<td>(\beta) (°)</td>
<td>106.79(6)</td>
<td>107.07(5)</td>
<td>107.04(3)</td>
<td>107.02(4)</td>
<td>107.02(3)</td>
</tr>
<tr>
<td>V (Å\textsuperscript{3})</td>
<td>382.5</td>
<td>381.6</td>
<td>379.8</td>
<td>379.6</td>
<td>379.1</td>
</tr>
</tbody>
</table>

These neutron data allow determination of the behaviour of the lattice parameters as a function of T (table 3). These results show a consistent increase in cell volume, although not all lattice parameters appear to change monotonically. Within error, there is virtually no change in \(\beta\), and changes in the other parameters are also relatively small compared to the errors.

The magnetic diffraction above and below \(T_N\) shows no qualitative difference, but rather a straightforward change in the magnitudes of the Bragg peaks. It appears that the ground state is the same above and below \(T_N\), so the cusp relates to a change in the glassy aspects of the behaviour rather than the equilibrium magnetic ordering.

Neutron diffraction measurements were repeated for all temperatures with the sample in an applied field of 1T. In all cases, the 1T and 0T measurements for a given temperature were indistinguishable. When it is recalled that the sample was allowed to relax before the diffraction measurements were undertaken, this supports the idea that, given time to relax, the field-cooled sample eventually falls into the state obtained on warming the sample, which appears to be the same as the state arising in zero field.

6. Conclusions

Three distinct magnetic phases have been shown for Fe\textsubscript{0.5}Ni\textsubscript{0.5}PS\textsubscript{3} one at low temperatures below \(T_N\) \(\sim\) 14K, a second at intermediate temperatures and paramagnetism above \(T_N\) \(\sim\) 100K. However, \(T_N\) is dependent on the direction of the ramping of the temperature. Relaxation has also been observed from the FCC state to that of the ZFCW/FCW states over a period of approximately two hours (time constant of \(\sim\) 48 minutes). Therefore, two of the three requirements for magnetic glass identification have been sufficiently satisfied to conclude that the intermediate phase is likely a magnetic glass, induced by mixed exchange and anisotropy.

Neutron diffraction measurements have shown that the magnetic structure of Fe\textsubscript{0.5}Ni\textsubscript{0.5}PS\textsubscript{3} is similar to that of CoPS\textsubscript{3} and NiPS\textsubscript{3}. Furthermore, from these measurements it was determined that the direction of the magnetic moments is unlikely to lie along the \(b\) axis. Magnetic susceptibility measurements have shown that the direction of the magnetic moments are unlikely to lie in the \(c^*\) direction either. Therefore, it is suggested that the moments lie along the \(a\) axis.
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References


